

Group D Nuclear Reactions - Theory

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International Workshop on Acceleration and Applications of Heavy Ions



Plan of the talk

1 Introduction

- Interactions of Nuclei: Scattering and Reaction
- Cross-Sections

2 Theoretical Models

- Elastic Scattering
- Optical Model
- Coupled Channels Model
- α cluster Model for Transfer

3 Data Analysis with Fresco

- Fresco and Sfresco
- Data Analysis



Interactions of Nuclei

- Scattering

- Elastic: Nuclei *and* energy remain the same

$$a + A \Rightarrow a + A, \quad Q = 0$$

$$(Q = (m_{\text{initial}} - m_{\text{final}})c^2)$$

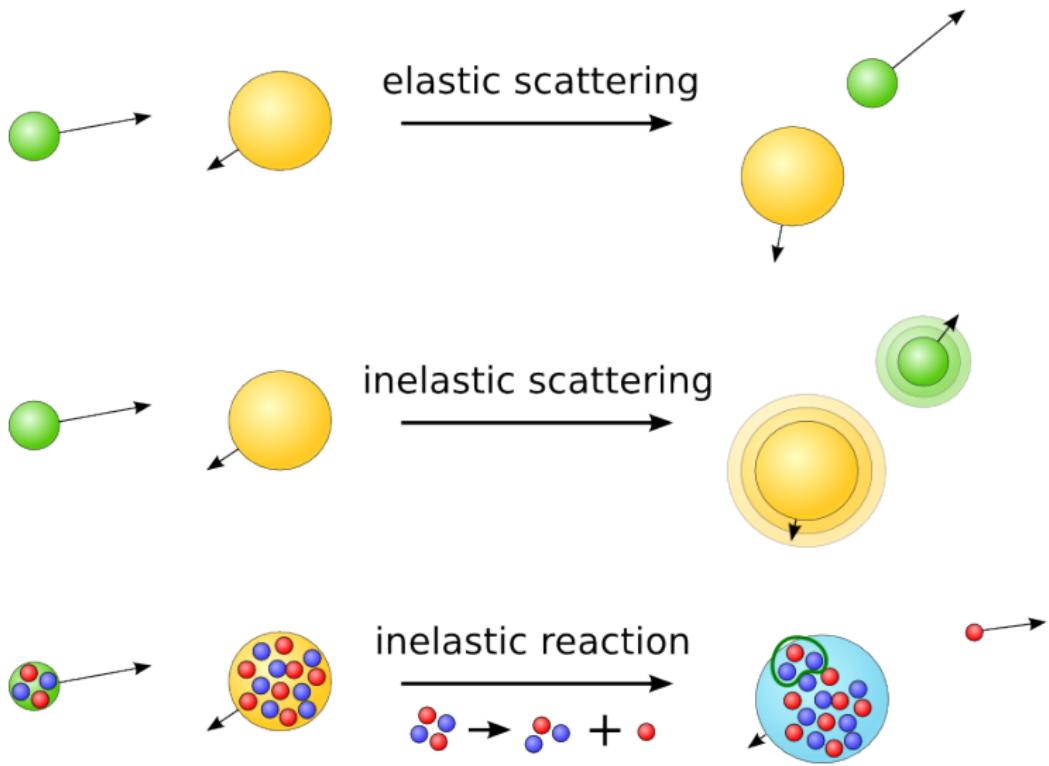
- Inelastic

$$a + A \Rightarrow a^* + A^*, \quad Q < 0$$

- Inelastic Reaction

$$a + A \Rightarrow b + B$$







Radiative capture

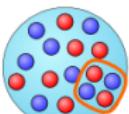
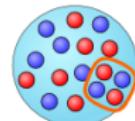
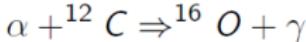
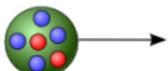
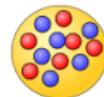
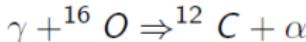
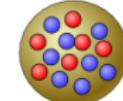
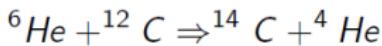


Photo reaction

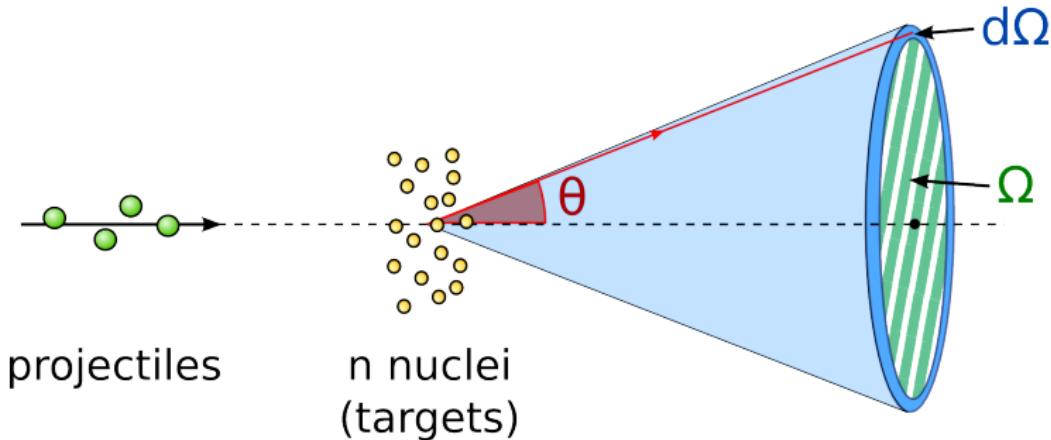


Transfer reaction



Observable: Cross-section

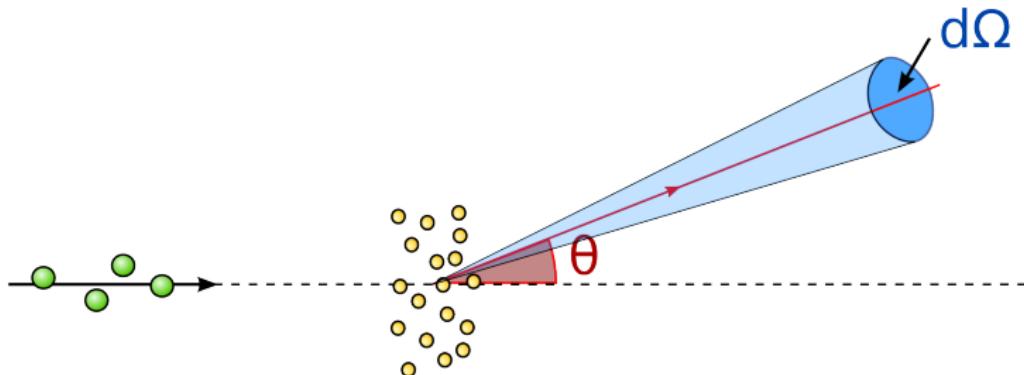
$$\sigma = \frac{\text{detected particles per unit time}}{\text{projectiles current} \times \text{count of target nuclei}}$$



Differential Cross-Section

$\frac{d\sigma}{d\Omega}$ - differential cross section - the cross section per unit of a solid angle.

$$\int_0^\pi \int_0^{2\pi} \frac{d\sigma}{d\Omega} \sin \theta d\varphi d\theta = \sigma$$



projectiles

n nuclei
(targets)



Theoretical Models

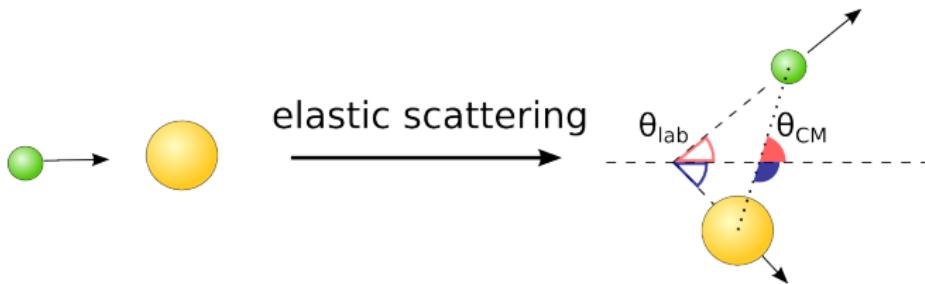
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Elastic Scattering



Schrödinger's equation:

$$\left[-\frac{\hbar^2}{2m_1} \nabla_1^2 - \frac{\hbar^2}{2m_2} \nabla_2^2 + V(\vec{r}_1, \vec{r}_2) \right] \psi = E\psi$$

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

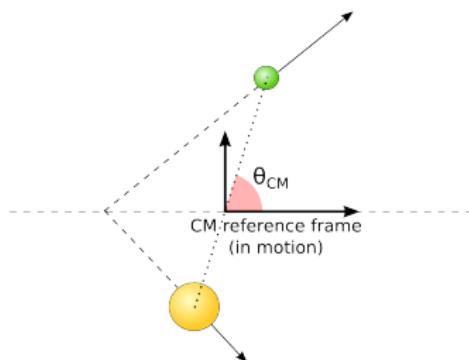
In the CM reference frame:

$$M = m_1 + m_2$$

$$\underbrace{-\frac{\hbar^2}{2M} \nabla_{R_{CM}}^2 \psi}_{= \varepsilon_0 \psi} - \underbrace{\left(-\frac{\hbar^2}{2\mu} \nabla_r^2 + V(r) \right) \psi}_{= \varepsilon \psi} = E\psi$$

$$\nabla^2 = \frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} + \frac{\partial^2}{\partial z_i^2}$$

Elastic Scattering - Reduction to One-Body Problem



Center of mass motion:

$$-\frac{\hbar^2}{2M} \nabla_{R_{CM}}^2 \psi = \varepsilon_0 \psi$$

Relative motion:

$$\left(-\frac{\hbar^2}{2\mu} \nabla_r^2 + V(r) \right) \psi = \varepsilon \psi$$

Potential of the Interaction

$$V_{\text{eff}} = V_C(r) + V_I(r) + V_N(r) + V_S$$

- Coulomb Potential (R_C is the Coulomb radius)

$$V_C(r) = \frac{Z_1 Z_2 e^2}{r}, \quad r > R_C$$

$$V_C(r) = \frac{Z_1 Z_2 e^2}{2R_C} \left(3 - \frac{r^2}{R_C^2} \right), \quad r < R_C$$

- Centrifugal Potential

$$V_I(r) = \frac{I(I+1)\hbar^2}{2\mu r^2}$$

- Nuclear Potential $\rightarrow ??$

Optical Model

Elastic scattering. Imaginary potential \rightarrow absorbtion.

Optical Potential:

$$V_N(r) = V(r) + iW(r)$$

- Phenomenological Potentials

- Woods-Saxon: $V(r) = \frac{V}{1 + e^{(r-R_0)/a}}$

$$R_0 = r_0(A_{\text{projectile}}^{1/3} + A_{\text{target}}^{1/3})$$

- Woods-Saxon Squared: $V(|r|) = \frac{V}{(1 + e^{(r-R_0)/a})^2}$

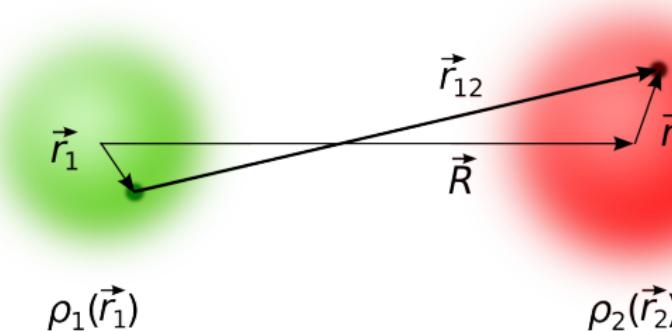
- Folding Potential

$$V(r) = \int \int \rho_1(\vec{r}_1) V \rho_2(\vec{r}_2) d^3\vec{r}_1 d^3\vec{r}_2$$

Here we can use different functions for ρ and V



Folding potentials



- Double-folding

$$V_{DF} = \int \int \rho_{\text{proj}}(\vec{r}_{\text{proj}}) V(\vec{r}_{12}) \rho_{\text{targ}}(\vec{r}_{\text{targ}}) d^3 \vec{r}_{\text{proj}} d^3 \vec{r}_{\text{targ}}$$

- Single-folding

$$V_{SF} = \int \rho_{\text{targ}}(\vec{r}_{\text{targ}}) V(\vec{r}_{12}) d^3 \vec{r}_{\text{targ}}$$

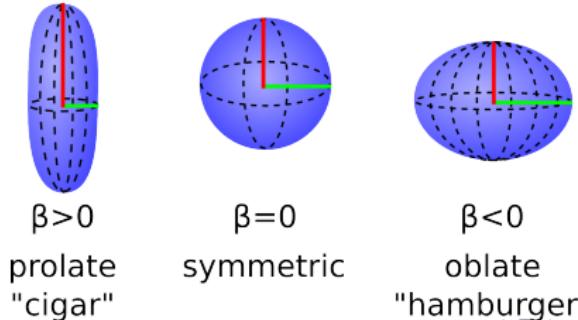
Coupled Channels

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V_{opt} + \frac{I(I+1)}{\hbar^2} - E \right] \psi_0 = \begin{cases} 0 \Rightarrow \text{Optical model} \\ U(r)\psi_1 \Rightarrow \text{Inelastic channel} \end{cases}$$

In the Coupled Channels Model, we deform the potential by modifying the radius:

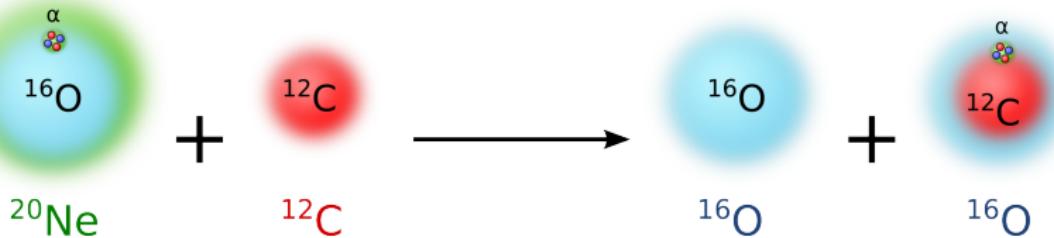
$$R \rightarrow r_0 A^{1/3} [1 + \beta Y_{20}(\theta, \varphi)]$$

So the potentials are no more spherically symmetrical.



α Cluster Transfer

It is known that Ne and C have α cluster structure ($N\alpha$)



$$V_N = V_{\alpha-\text{O}} + V_{\alpha-\text{C}} + V_{\text{O-C}}$$

Data Analysis with Fresco

1 Introduction

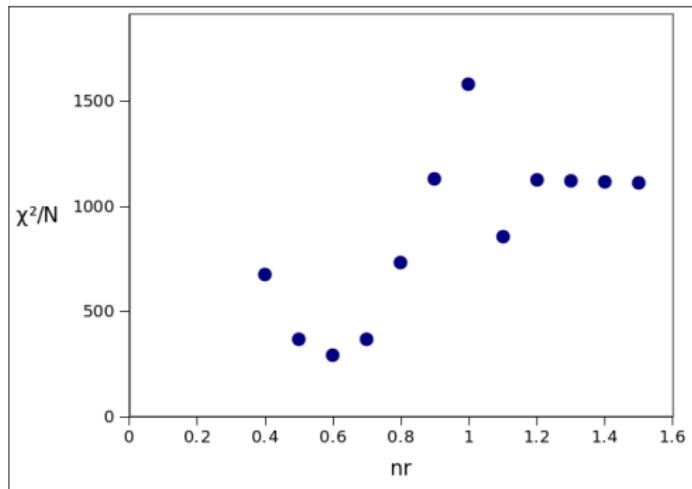
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Fresco and Searching Fresco (sfresco)

- We used the program **fresco** to do the calculations.
- **sfresco** is the searching version of **fresco**. It tries to optimize the parameters for best fitting.



Analyzing Experimental Data

We try to describe the data obtained by Group C

- Optical Model
- Coupled Channels Model
- Alpha Clusters Model



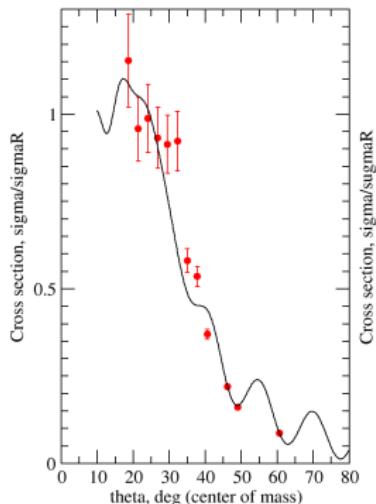
Optical Model - Woods-Saxon Potential

Woods-Saxon Squared for the imaginary part

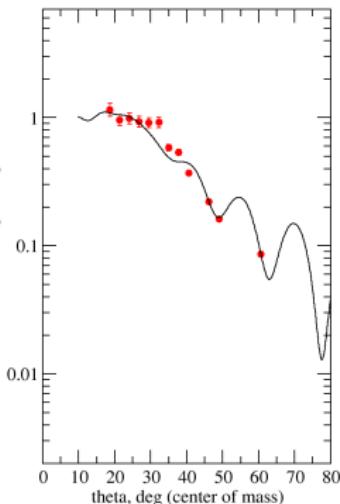
$$V(r) = \frac{V}{1 + e^{(r-R_0)/a_0}}$$

$$W(r) = \frac{W}{(1 + e^{(r-R_i)/a_i})^2}$$

Linear scale



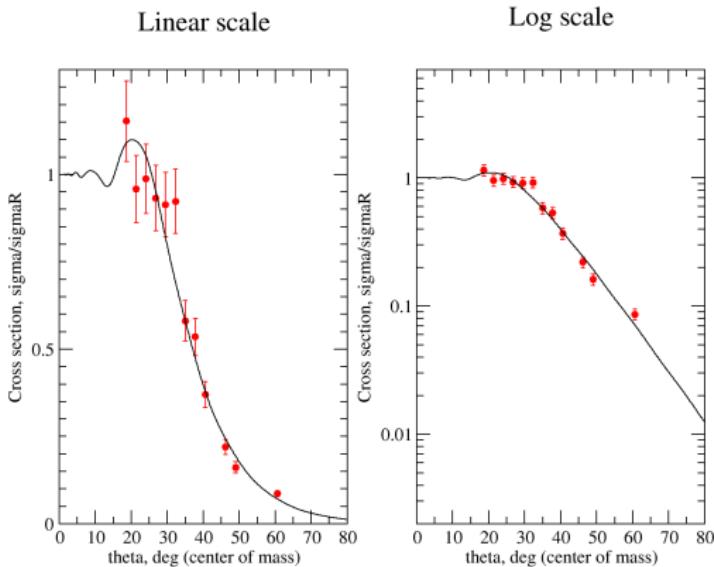
Log scale



$$\begin{aligned} V &= 146.2 \text{ MeV}; r_0 = 1.30 \text{ fm}; a_0 = 0.637 \text{ fm}; \\ W &= 5.77 \text{ MeV}; r_i = 1.28 \text{ fm}; a_i = 0.648 \text{ fm} \end{aligned}$$

Optical Model - Double-Folding Potential

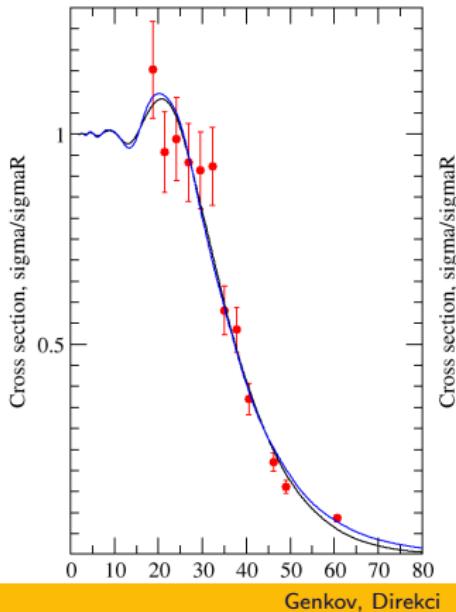
M3Y double folding potential, Gaussian densities



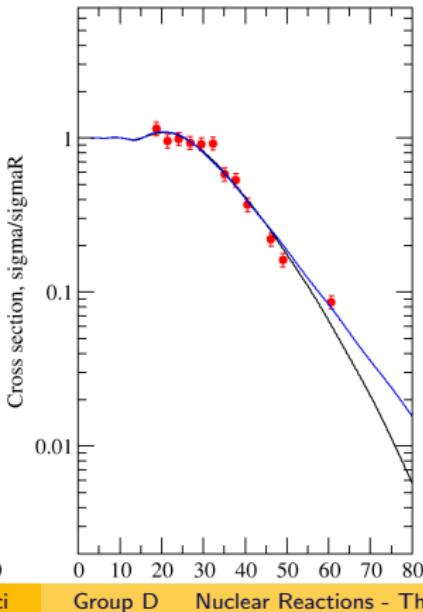
$$N_R = 0.283, \quad N_i = 0.321$$

Double-Folding Potential: Optical vs. Coupled Channels Models

Linear scale

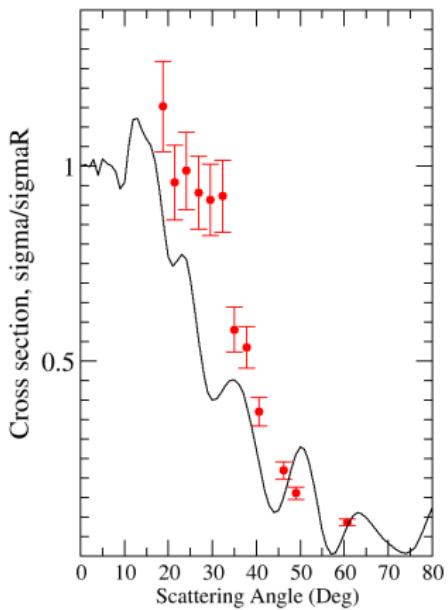


Log Scale

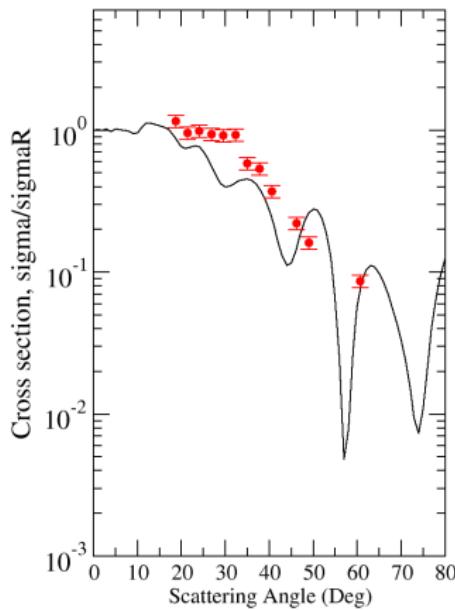


α Cluster Transfer

Linear scale



Log scale



Conclusion

- Optical model with phenomenological and double folding potentials are successful to describe the experimental data.
- We need a deep real potential ($V \approx 150$ MeV) and a shallow imaginary potential ($W \approx 5$ MeV).
- These potential values prove that the interaction of projectile and target nuclei take place at the surface.
- α -transfer should be studied further in the future to understand the following difference in the total cross-sections:

σ_{total}	σ_{CC}	σ_{OM}	$\sigma_{\alpha\text{trans}}$
1142.20 mb	1060.93 mb	1130.82 mb	1583.13 mb

Thanks

People who helped us

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- **ALL OF YOU!**

